

Graph Algorithm for the Simulation of the Interaction between Particles

Stéphane Després, Roland Hildebrand, Aude Maignan

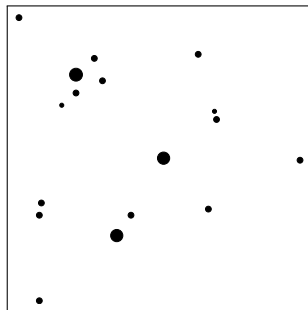
September 24 2012

Dynamical system of interacting particles

- Mass m_i
- Position $\chi_i(t)$
- Momentum $\pi_i(t) = m_i \dot{\chi}_i(t)$

Newton's second law of motion:

$$m_i \ddot{\chi}_i = G \sum_{k \neq i} \frac{m_i m_k (\chi_k - \chi_i)}{|\chi_k - \chi_i|^3}$$



Hamiltonian equation

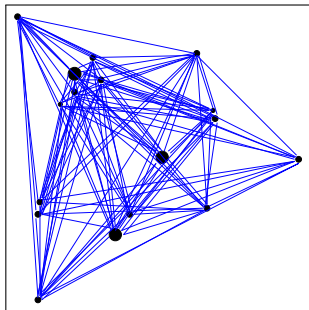
$$H(\chi, \pi) = \sum_i \frac{\|\pi_i\|^2}{2m_i} - \sum_{i < k} G \frac{m_i m_k}{|\chi_i - \chi_k|}$$

Equation of motion

$$\begin{cases} \frac{d\chi}{dt} = \frac{\partial H}{\partial \pi} \\ \frac{d\pi}{dt} = -\frac{\partial H}{\partial \chi} \end{cases} .$$

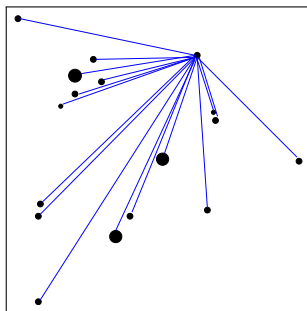
Problems of direct integration

- time computation
- stability
- scale problems



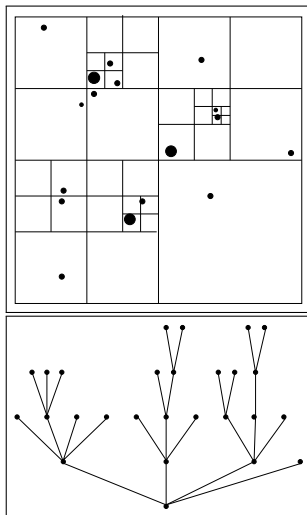
N -body simulation

- Evaluating the forces on individual particles
- Integrating the orbits of particles



Evaluating the forces on individual particles : Tree methods

- Barnes and Hut[1986,1989]
- Fast multipole method [Carrier, Greengard, 1988]



Integrating the orbit of particles

- Kepler laws for $n = 2$
- Runge-Kutta with adaptive time step
- Symplectic integration [Hierarchical Jacobi Symplectic, Beust 2003]

Integrating the orbit of particles

- Kepler laws for $n = 2$
- Runge-Kutta with adaptive time step
- Symplectic integration [Hierarchical Jacobi Symplectic, Beust 2003]
 - the stellar system is static and modelled with a binary hierarchical tree.
 - $H = H_A + R$ with $R \ll H_A$
 - H_A : The sum of Keplerian Hamiltonians
 - R : the residual mutual perturbations

The Classical N -body problem

The Nbodies Softwares

Sverre Aarseth's Institute of Astronomy Homepage <http://www.astr.cam.ac.uk/~sverra/>

Sverre Aarseth

AR Round N-Body Day, Mountairmer and WML Life Achievement

Home News **N-Body & Simulations** Publications Interviews Contact Us

Welcome to Sverre Aarseth's Home Page!





Sverre's web pages are hosted by the University of Cambridge Institute of Astronomy



Original site design by www.star.ucl.ac.uk/~m1000/

1 of 1 06/09/2012 16:20


Sverre Aarseth's Institute of Astronomy N-Body and Dow... <http://www.astr.cam.ac.uk/~sverra/web/pages/nbody.htm>

Sverre Aarseth

AR Round N-Body Day, Mountairmer and WML Life Achievement

Home News **N-Body & Simulations** Publications Interviews Contact Us

N-Body



Introduction

My research into numerical simulations of many-body (N-body) gravitational interactions spans over 40 years, and is reported in several [publications](#) and my book. I have developed a set of FORTRAN codes which describe the dynamics very closely, and these are regarded by many as the de facto standard. I have explained below how you can download these codes for execution on your local computer system.

Downloads

For those of you with an FTP tool, you can access my public N-body directory at <http://www.astr.cam.ac.uk/~sverra/>. Your browser (such as Microsoft's Internet Explorer) cannot present you with an FTP viewer window when you click on this link. If so, you should be able to copy whole code group directories at a time.

Otherwise please use the links below to identify and download the specific files that you need.

At any time if you'd like further information, please do [contact](#).

Code Group	Filename	Filename Size	Notes
	CODES	2	General textual description of the codes.
nbody1	SCALAPACK	1	The version nbody1 tar 2 is a full SCALAPACK
	nbody1.tar.gz	26	This is the nbody1.tar 2 in the tar version with mixed SCALAPACK/4x4 precision which is kept for historical reasons.
	nbody1.tar.gz.2	26	This is the nbody1.tar 2 in the tar version with mixed SCALAPACK/4x4 precision which is kept for historical reasons.
	nbody1.tar.gz	140	The version nbody1 tar 2 is a full SCALAPACK
nbody2	nbody2.tar.gz	213	Directory nbody2 contains the code NBODY2 and an article describing it. The whole file is called nbody2.tar.gz. Inside is the code. The article should be read and sent to sverra@astr.cam.ac.uk .
	nbody2.tar.gz.2	185	The version nbody2 tar 2 is the tar version with mixed SCALAPACK/4x4 precision which is kept for historical reasons.
	nbody2.tar.gz	204	One force polynomial with full regularization.
nbody4	SCALAPACK	1	The public SCALAPACK or micro-Grange code NBODY4
	nbody4.tar.gz	340	and the equivalent force finite code for workstations or laptops called BRUTA
	nbody4.tar.gz.2	240	The current versions were updated 13 March 2006.
	user_nbody4.pdf	1	Introduction to running simulations
nbody6	SCALAPACK	1	The code NBODY6 has been heavily used but
	nbody6.tar.gz	258	is no longer updated. NBODY6 is recommended instead
nbody8	SCALAPACK	1	SCALAPACK contains some past history of NBODY8.
	nbody8.tar.gz	640	Further practical info can be found in directory docs.
	nbody8.tar.gz.2	440	nbody8 tar gz holds the new standard code NBODY8.
	nbody8.pdf	200	The previous 2000 version has been replaced by the code tar.gz.
	nbody8.pdf.2	200	The previous 2000 version has been replaced by the code tar.gz.
	user_nbody8.pdf	1	Reader Squamant's corresponding parallel code, nbody8n, can be obtained here.
	user_nbody8.pdf	1	Introduction to running simulations.
	nbody8n.pdf	1	SCALAPACK contains some past history of NBODY8.
	nbody8n.pdf.2	1	SCALAPACK contains some past history of NBODY8.
	nbody8n.pdf	45	SCALAPACK contains some past history of NBODY8.

1 of 2 07/09/2012 11:24

Outlines

- Tree model of the hierarchical partition
 - ⇒ extract more information than Barns and Hut
- change of coordinates and symplectic integration
 - ⇒ extension of Beust method
- Algorithm and Software

The logo for GASIP, with the letters in a bold, metallic, 3D font. A bright blue starburst effect is positioned above the letter 'A'.

Graph model

Complete graph \mathcal{G} of interactions

- Node i : Particle
- Weighted edge (i, j) : $w_{ij} = G \frac{m_i m_j}{|x_i - x_j|}$

Partition of \mathcal{G} of parameter ϵ . If A and B are 2 parts of the partition :

- $i \in A$ and $j \in B \implies w_{ij} \leq \epsilon$
- $i \in A$ and $j \in A \implies$ there exists a path from i to j such that every edge weight are greater than ϵ

Complexity of the hierarchical partition algorithm: $O(N \log N)$

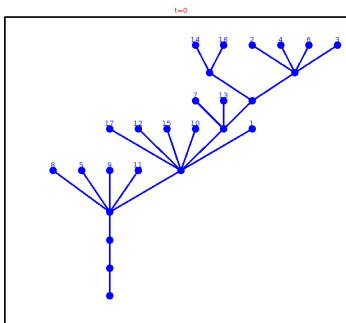
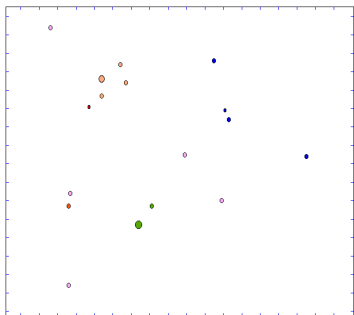
The adjacency matrix of \mathcal{G} : $A = \begin{cases} a_{ij} = \frac{m_i m_j}{\|x_i - x_j\|} & \text{if } i \neq j \\ -\sum_{j, j \neq i} a_{ij} & \text{else} \end{cases}$

The partition is obtained by :

- Comparing at most one time the coefficients of A with the value of ϵ
- Executing rows and columns permutations.

A is equivalent to $D_A = \begin{pmatrix} \boxed{B_1(\epsilon)} & & & & \\ & \leq \epsilon & & & \\ & & \cdot & & \\ & & & \cdot & \\ & & & & \leq \epsilon \\ & & & & & \boxed{B_k(\epsilon)} \end{pmatrix}$

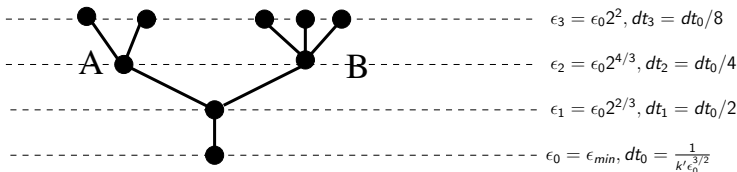
Partition tree



- The set of interactions $particle_A \leftrightarrow particle_B$ are approximated by an interaction $center_mass_A \leftrightarrow center_mass_B$
- The time step of integration depends on the level in the tree.

ϵ is multiplied by $2^{2/3}$ at each step \implies

the time step of integration is divided by 2



ϵ lower bound

Theorem

- Let λ_2 be the second eigenvalue of the Laplacian matrix

$$L = \begin{cases} l_{ij} = -w_{ij} & \text{if } i \neq j \\ -\sum_{j, j \neq i} l_{ij} & \text{else} \end{cases}$$

- let n the size of the matrix L .

Then for all

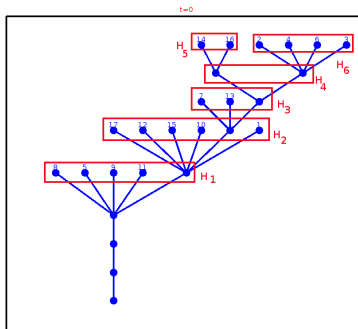
$$\epsilon < \frac{\lambda_2}{n}$$

the partition is reduced to the trivial partition.

Symplectic integration

$$H = H_A + R$$

- $H_A = H_1 + H_2 + H_3 + H_4 + H_5 + H_6$ contains the main interactions
- R contains all the residual interactions



Coordinates change is performed in order to incorporate the tree structure explicitly.

Extended Basis

Coordinate change $(\chi_i, \pi_i)_{i \in \{1..n\}} \rightarrow (x_i, p_i)_{i \in \{1..N\}}$ with

- n : number of particles
- N : number of nodes

Coordinate change Formulas

$$x = A\chi \text{ and } p = C\pi$$

with $A \in \mathcal{M}_{3N \times 3n}$ and $C \in \mathcal{M}_{3N \times 3n}$ depend on

- the masses
- the structure of the tree.

Extended coordinates is a generalization of the base of Jacobi

The original Hamiltonian

$$H(\chi, \pi) = \pi^T \Delta \pi + f(\chi), \quad (1)$$

in the new coordinates, H has a similar expression

$$H(x, p) = p^T D p + F(x). \quad (2)$$

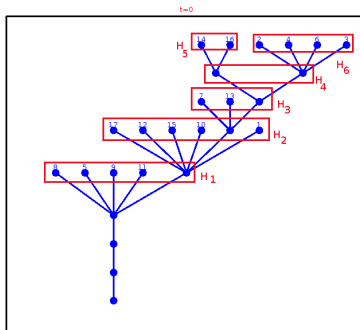
Theorem

- Let $(\chi(t), \pi(t))$ be a trajectory of system (1). Then $(x(t), p(t)) = (A\chi(t), C\pi(t))$ is a trajectory of system (2).
- Let $(x(t), p(t))$ be a trajectory of system (2). Then $(\chi(t), \pi(t)) = (C^T x(t), A^T p(t))$ is a trajectory of system (1).

Integration of $H(x, p) = \sum H_i(x, p) + R(x)$

$$H(x, p) = \sum H_i(x, p) + R(x)$$

- The H_i are integrated separately in the (x, p) coordinates.
 - If H_i contains 2 particles → Kepler method
 - else Runge Kutta is used with a time step determined by the tree.
- R is integrated.



Symplectic integration part of

$$H(x, p) = \sum H_i(x, p) + R(x) \text{ in } [t, t + \tau]$$

- 1 Integration of $\sum H_i(x, p)$ at a time step $\tau/2$. We obtain
$$(\tilde{x}(t + \tau/2), \tilde{p}(t + \tau/2)) = (x(t), p(t)) + \int_t^{t+\tau/2} \left(\frac{\partial \sum H_i}{\partial p}, -\frac{\partial \sum H_i}{\partial x} \right)$$
- 2 Integration of R at a time step τ . We obtain
$$(x', p') = (\tilde{x}(t + \tau/2), \tilde{p}(t + \tau/2)) + \tau \frac{\partial R}{\partial x}(\tilde{x}(t + \tau/2))$$
- 3 Integration of $\sum H_i$ at a time step $\tau/2$. Finally,
$$(\tilde{x}(t + \tau), \tilde{p}(t + \tau)) = (x', p') + \int_{t+\tau/2}^{t+\tau} \left(\frac{\partial \sum H_i}{\partial p}, -\frac{\partial \sum H_i}{\partial x} \right).$$

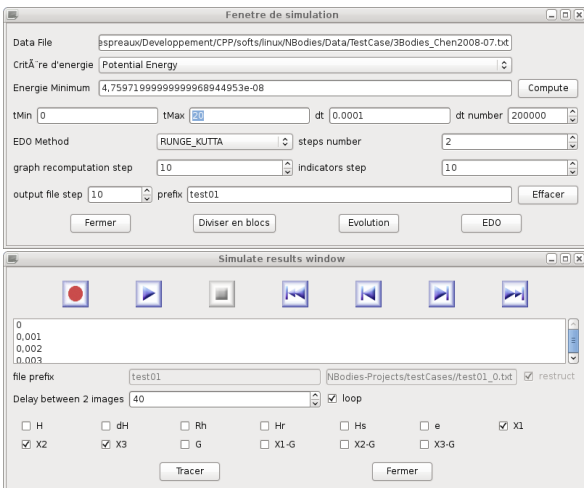
$(\tilde{x}(t + \tau), \tilde{p}(t + \tau))$ is the approximation of $(x(t + \tau), p(t + \tau))$.

Algorithm

- ① $t = t_{min}$
- ② While ($t < tMax$) do
 - Compute the tree
 - Perform the coordinate transformation $(\chi, \pi) \rightarrow (x, p)$
 - Construct the individual parts of the Hamiltonian in the new coordinate system:
$$H(x, p) = \sum H_i(x, p) + R(x)$$
 - $n = 0$
 - While ($n < nTreeSteps$) and ($t < tMax$) do
 - Integrate $(\dot{x}, \dot{p}) = (\frac{\partial H}{\partial p}, -\frac{\partial H}{\partial x})$ with the combined integration method and a bottom time step of $\tau = \frac{1}{k\epsilon^{3/2}}$.
 - $n = n + 1$
 - $t = t + \tau$
 - Perform the reverse coordinate transformation:
$$(x, p) \rightarrow (\chi = C^T x, \pi = A^T p)$$



GASIP



Fenetre de simulation

Data File: `gspreaux/Developpement/CPP/softs/linux/NBodies/Data/TestCase/3Bodies_Chen2008-07.txt`

Critère d'énergie: Potential Energy

Energie Minimum: `4,75971999999999968944953e-08` Compute

tMin: 0 tMax: 20 dt: 0.0001 dt number: 200000

EDO Method: RUNGE_KUTTA steps number: 2

graph recomputation step: 10 indicators step: 10

output file step: 10 prefix: test01 Effacer

Fermer Diviser en blocs Evolution EDO

Simulate results window

0
0.001
0.002
0.003

file prefix: test01 `NBodies-Projects/testCases/test01_0.txt` restruct

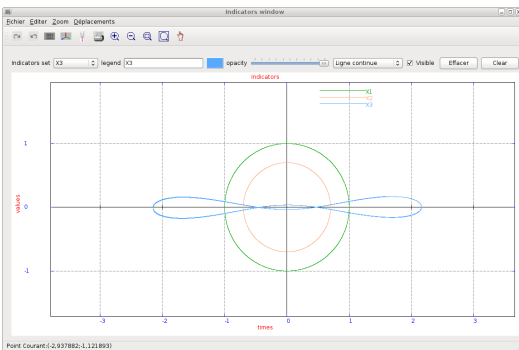
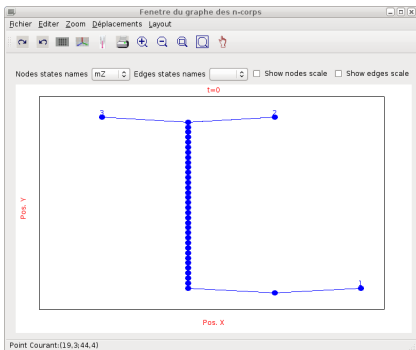
Delay between 2 images: 40 loop

H dH Rh Hr Hs e X1
 X2 X3 G X1-G X2-G X3-G

Tracer Fermer



GASIP



Conclusion

Software



- Dynamical partition tree
 - ⇒ Flocking simulation
 - ⇒ Wireless sensor networks
- New augmented coordinates system generalized the base of Jacobi.